

Solutions

(A1) A numerical scheme is called convergent if for any fixed point  $(x^*, t^*)$  in the computational domain we have:

From

(1)  $\Delta t \rightarrow 0$ ,  $\Delta x \rightarrow 0$  and  $\mu = \frac{\Delta t}{(\Delta x)^2} = \text{const}$

(2)  $x_j \rightarrow x^*$  and  $t_n \rightarrow t^*$

it follows that

$$U_j^n \rightarrow u(x^*, t^*)$$

1 point

(A2)  $\checkmark$  let  $e_j^n = U_j^n - u(x_j, t_n)$  be (local) error in the numerical scheme

Then  $E^n = \max \{ |e_j^n|, j = \underbrace{0, 1, 2, \dots, J}_{\text{all mesh points}} \}$

0.5 point

Let  $\bar{T} = \checkmark$  upper bound for the truncation error,

i.e.  $|T_j^n| \leq \bar{T}$ . Let us set  $\bar{T} = C \Delta t$

0.5 point

An estimate holds:

$$E^n \leq n \bar{T} \Delta t = \underbrace{(n \Delta t = t_F)}$$

therefore:

$$= \bar{T} t_F = C \Delta t t_F$$

1 point

(since CFL number is

fixed, so is C)

(2)

(A3) Rewrite the given scheme:

$$U_j^{n+1} - U_j^n = \frac{\Delta t}{(\Delta x)^2} \left[ D_{j+1/2} (U_{j+1}^n - U_j^n) - D_{j-1/2} (U_j^n - U_{j-1}^n) \right] =$$

$$= \frac{\Delta t}{(\Delta x)^2} \left[ D_{j+1/2} U_{j+1}^n - (D_{j+1/2} + D_{j-1/2}) U_j^n + D_{j-1/2} U_{j-1}^n \right]$$

with  $\mu = \frac{\Delta t}{(\Delta x)^2}$ :

$$U_j^{n+1} = \mu D_{j-1/2} U_{j-1}^n + \underbrace{\left( 1 - \mu (D_{j-1/2} + D_{j+1/2}) \right)}_{\text{denote by } b_j} U_j^n + \mu D_{j+1/2} U_{j+1}^n \quad 1 \text{ pt.}$$

 $\mu D_{j \pm 1/2}$  can be seen as a local CFL number

$$\left\{ \begin{array}{l} \mu D_{j \pm 1/2} > 0 \\ b_j^n > 0 \text{ provided the CFL number is small enough} \end{array} \right\} \quad 1 \text{ pt.}$$

positivity test

sum-property test:

$$\mu D_{j-1/2} + b_j + \mu D_{j+1/2} = 1 \quad 1 \text{ pt}$$

(B1) CFL condition is satisfied if the domain of dependence of the equation lies within the domain of dependence of the numerical scheme. 1 pt.

The CFL condition can be written as a bound imposed on the CFL number, for this case

$$\gamma = a \frac{\Delta t}{\Delta x}$$

(B2) Truncation error is the residual obtained when the exact solution of the PDE is substituted into the numerical scheme. 1 pt

Let  $u(x, t)$  be exact solution of the PDE

$$u_t + a u_x = 0.$$

Consider  $u_j^{n+1}$  and  $u_j^n$  and omit the superindices  $n$  and  $n+1$ :

(3)

(derivative with respect to the space variable  $x$ )

$$u_j = u_{j+1/2} + u'_{j+1/2} \frac{-h}{2} + u''_{j+1/2} \frac{h^2}{4} - u'''_{j+1/2} \frac{h^3}{6} + O(h^4)$$

with  $h = \Delta x$

Similarly for  $u_{j+1}^n$  and  $u_{j+1}^{n+1}$ :

$$u_{j+1} = u_{j+1/2} + u'_{j+1/2} \frac{h}{2} + u''_{j+1/2} \frac{h^2}{4} + u'''_{j+1/2} \frac{h^3}{6} + O(h^4)$$

Hence:

$$u_j + u_{j+1} = 2u_{j+1/2} + 2u''_{j+1/2} \frac{h^2}{4} + O(h^4) \quad \text{for } t = t_n \text{ and } t = t_{n+1}$$

Therefore

$$(u_j^{n+1} + u_{j+1}^{n+1}) - (u_j^n + u_{j+1}^n) = 2u_{j+1/2}^{n+1} + 2(u''_{j+1/2})^{n+1} \frac{h^2}{4} + O(h^4) - 2u_{j+1/2}^n - 2(u''_{j+1/2})^n \frac{h^2}{4} + O(h^4)$$

and

$$\frac{(u_j^{n+1} + u_{j+1}^{n+1}) - (u_j^n + u_{j+1}^n)}{2\Delta t} = \frac{u_{j+1/2}^{n+1} - u_{j+1/2}^n}{\Delta t} + \frac{(u''_{j+1/2})^{n+1} - (u''_{j+1/2})^n}{\Delta t} \cdot \frac{h^2}{4} + O(h^4)$$

(derivative with respect to  $t$ )

$$= (\ddot{u} + O(\Delta t)^2)_{j+1/2}^{n+1/2} + (\ddot{u} + O(\Delta t)^2)_{j+1/2}^{n+1/2} + O(h^4) \frac{1}{\Delta t}$$

Due to the symmetry in the space and time of the two fractions in formula (3), we have

$\left(\frac{h^2}{4}\right)$  [1 point]

$$\frac{(u_{j+1}^n + u_{j+1}^{n+1}) - (u_j^n + u_j^{n+1})}{2\Delta x} = \underbrace{(u' + O(\Delta x)^2)_{j+1/2}^{n+1/2}}_{O(\Delta t)^2} + \underbrace{(\ddot{u}' + O(\Delta x)^2)_{j+1/2}^{n+1/2} \frac{(\Delta t)^2}{4}}_{O(\Delta t)^2} + \underbrace{\frac{1}{\Delta x} O(\Delta t)^3}_{O(\Delta t)^2}$$

Adding up the two expressions, we get

if CFL number is bounded

$$\frac{\dots}{2\Delta t} + a \frac{\dots}{2\Delta x} = \underbrace{(\ddot{u} + a u')_{j+1/2}^{n+1/2}}_{=0 \text{ because } u(x,t) \text{ is the solution}} + O(\Delta x)^2 + O(\Delta t)^2 \equiv$$

$$\equiv O(\Delta x)^2 + O(\Delta t)^2 \leftarrow \begin{array}{l} \text{second order} \\ \text{in space and} \\ \text{time.} \end{array}$$

[1 point]

(4)

B3) damping error =  $|\lambda| - 1$  ← exact damping = no damping

phase error =  $\arg(\lambda) - (-\sqrt{\xi})$   
↑  
 exact shift in phase

Thus, to evaluate the errors we need to estimate  
 $|\lambda|$  and  $\arg(\lambda)$

1 point

$$\lambda = 1 - \sqrt{\xi} + \sqrt{\xi} \cdot e^{-i\xi} = \sqrt{\xi} (\cos \xi - i \sin \xi) = 1 + \sqrt{\xi} \cos \xi - i \sqrt{\xi} \sin \xi$$

$$|\lambda|^2 = (1 - \sqrt{\xi} + \sqrt{\xi} \cos \xi)^2 + \xi \sin^2 \xi =$$

$$= (1 - \sqrt{\xi})^2 + 2\sqrt{\xi}(1 - \sqrt{\xi}) \cos \xi + \xi (\sin^2 \xi + \cos^2 \xi) =$$

0.5 point

$$= 1 - 2\sqrt{\xi} + \xi + 2\sqrt{\xi}(1 - \sqrt{\xi}) \cos \xi + \xi =$$

$$1 - 2\sqrt{\xi} + 2\xi + 2\sqrt{\xi}(1 - \sqrt{\xi}) \left(1 - 2 \sin^2 \frac{\xi}{2}\right) =$$

$$= 1 - 2\sqrt{\xi}(1 - \sqrt{\xi}) + 2\sqrt{\xi}(1 - \sqrt{\xi}) \left(1 - 2 \sin^2 \frac{\xi}{2}\right) =$$

$$= 1 + 2\sqrt{\xi}(1 - \sqrt{\xi}) \left[-1 + 1 - 2 \sin^2 \frac{\xi}{2}\right] =$$

$$= 1 + 2\sqrt{\xi}(1 - \sqrt{\xi})(-2) \sin^2 \frac{\xi}{2} = 1 - 4\sqrt{\xi}(1 - \sqrt{\xi}) \sin^2 \frac{\xi}{2}$$

$$= 1 - O(\xi^2)$$

0.5 point

$$\Rightarrow \text{damping error} = O(\xi^2), \quad p=2.$$

$$\text{damping error} = |\lambda| - 1 = \sqrt{1 - O(\xi^2)} - 1 \approx 1 - O(\xi^2) - 1 = -O(\xi^2) = O(\xi^2)$$